Experimental Observations and Numerical Modeling of Coupled Microbial and Transport Processes in Variably Saturated Sand


ABSTRACT

An experimental and numerical investigation was conducted to study interactions between microbial dynamics and transport processes in variably saturated porous media. Experiments were conducted with constant, surface-applied water fluxes in duplicate, variably saturated, sand-filled columns that were uniformly inoculated with the bacterium *Pseudomonas fluorescens* HK44. The permeability of the sand in the columns was reduced by a factor of 45 during 1 wk of growth on glucose. Pressure heads increased (became less negative) at all measured depths, but significant increases in the apparent volumetric water contents were observed in only the upper 5 cm of the columns, corresponding to the areas with the highest concentrations of attached bacteria. A numerical model was used to simulate the experiments. The model accounted for the processes of water flow, solute and bacterial transport, cell growth and accumulation, glucose and O2 consumption, and gas diffusion and exchange. Observed changes in water content and pressure head were reproduced approximately using fluid-media scaling to account for an apparent surface-tension lowering effect. Reasonable correspondence was obtained between observed and simulated effluent data and final attached biomass concentration distributions using first-order reversible cell attachment and detachment kinetics. The attachment rate coefficients were based on particle-filtration theory and time-dependent detachment rate coefficients. The results of this study illustrate the potential importance of using fully coupled multifluid flow and multicomponent reactive transport equations to model coupled biogeochemical and transport processes in soils.

A significant body of research exists on the transport and biodegradation of various contaminants in saturated porous media systems that are representative of aquifer materials. Comparatively little work has been done to study biodegradation or bioremediation on a mechanistic basis in more complicated unsaturated porous media systems such as soils. Several studies have been conducted on the transport of bacteria in soils or unsaturated porous media under nongrowth conditions (Tan et al., 1992; Wan et al., 1994; Schäfer et al., 1998; Powelson and Mills, 1998; Jewett et al., 1999). A number of studies have also been conducted on the transport and degradation of model contaminants in soils (Estrella et al., 1993; Allen-King et al. (1996a, 1996b); Langner et al., 1998). However, the majority of the studies on contaminant transport and degradation in soils have represented biodegradation as a first-order decay process, without actually considering microbial processes and their interactions with other transport processes (Rockhold et al., 2004).

In the work of Estrella et al. (1993), the transport and fate of 2,4-dichloro-phenoxyacetic acid (2,4-D) was studied in both saturated and unsaturated soils. They determined that sorption had a slight but significant effect on 2,4-D transport, but that biodegradation was extensive, under both saturated and unsaturated conditions, following a lag period of up to 3 d. One of their most significant findings was that degradation rate parameters determined from batch experiments were significantly different from those determined in column experiments. The differences were attributed, in part, to possible O2 limitations in their saturated experiments that were either not present or less severe in their unsaturated experiments. Their column experiments were conducted under steady flow conditions, in both saturated and unsaturated soil, and they modeled these experiments using analytical and numerical models for aqueous phase transport of 2,4-D only. Other processes such as O2 transport in the aqueous and gas phases, interphase exchange, O2-limited 2,4-D degradation kinetics, and bacterial growth were not considered.

Langner et al. (1998) studied 2,4-D transport and degradation in batch and unsaturated column experiments with columns of different lengths. They determined that a single set of independently determined rate parameters from batch experiments could not describe 2,4-D degradation for all transport conditions. Apparent first-order degradation rate constants obtained from column data were found to be independent of column residence time, but increased with decreasing pore water velocity, especially at low velocities. They speculated that variations in apparent degradation rates were due to changes in microbial attachment and distribution under different flow rates, differences in the residence time or “local opportunity” time for degradation, or differences in nutrient desorption rates from the solid phase at different pore water velocities. They stated that there is a “need to develop a better understanding of coupled processes involving contaminant degradation and transport.”

The general objective of our work was to gain a better understanding of possible interactions and feedback mechanisms between bacterial growth, water flow, solute transport, and gas exchange in soils. These interactions have important consequences for applications such as

Abbreviations: DO, dissolved oxygen; EM, electromagnetic; MMS, minimal mineral salts; TDR, time domain reflectometry; UV, ultraviolet; 2,4-D, 2,4-dichloro-phenoxyacetic acid.
water and wastewater treatment or bioremediation of contaminated soils and aquifer sediments, as well as being of general ecological interest. Specific objectives of this study were to (i) quantify the impact of bacterial growth on the hydraulic properties of variably saturated sand and (ii) develop a numerical model to simulate these processes.

MATERIALS AND METHODS

Instrumentation

Duplicate, segmented, 7.62-cm-diameter (3 in. o.d.) acrylic columns were used in the experiments. The segments were machined in two heights, 1.8 and 6.0 cm, with finished inside diameters of 4.28 cm. The shorter segments were used in the upper, unsaturated part of the columns to provide higher spatial resolution in this region. A brass screen 0.149 mm (100 mesh) was placed in the bottom end-cap to retain the porous media. A relatively coarse screen was used instead of a finer, porous ceramic disc or stainless-steel plate, on which a negative pressure could be applied, to avoid clogging by bacterial cell aggregates or biofilms. The effluent lines were split so that samples could be collected periodically from one end for measurement of dissolved oxygen (DO). The other end was used for effluent outflow and head control. The column segments were held together with three threaded stainless-steel rods. The total height of the columns in which sand could be packed was approximately 55 cm. Experiments were conducted with sand packed to a height of 48 cm to allow room for drip emitters on the tops of the columns. The experimental setup is depicted in Fig. 1.

Alternating smaller segments in the columns were machined with inner water chambers and porous ceramic ring inserts, sealed in place with epoxy, to form tensiometers. The ceramic rings were cut from larger cylindrical ceramic tubes (Osmonics, Inc., Minnetonka, MN) with a nominal pore diameter of 10 μm, corresponding to a bubbling pressure of approximately 300 cm. Pressure transducers (Honeywell Micro Switch model 26PCBFA1G) were fitted to the tensiometers and were excited and monitored by a personal computer (PC) with a 12-bit data acquisition card (AT-MIO-64E-3), controlled by LabVIEW software (National Instruments Corp., Austin, TX).

The other set of the alternating smaller column segments and two of the larger segments in each column were machined and fitted with miniature, two-wire, time-domain-reflectometry (TDR) probes. The TDR probes were made from 0.9-mm-diameter stainless-steel wire plated with 14K gold to facilitate soldering to commercially available SMA bulkhead connectors (ITT Type 50-645-4524-310). The TDR probes were connected to a Tektronix Model 11801 digital sampling oscilloscope with an SD-24 sampling head (Tektronix, Beaverton, OR). The 20-GHz oscilloscope generates a step electromagnetic pulse with a rise time of 25 ps at the connector on the instrument (Kelly et al., 1995). The TDR signals were transferred from the oscilloscope to a PC via a general purpose interface bus that was controlled by a custom developed Microsoft Excel spreadsheet macro (VBA program). The spreadsheet program automatically plotted TDR waveforms and computed travel times. Apparent liquid saturations were calculated from travel times using regression equations established in independent calibration experiments.

Steady unsaturated fluxes of an autoclaved minimal mineral salts (MMS) solution containing 250 mg L⁻¹ of glucose were applied to the upper surfaces of the columns through cylindrical acrylic manifolds that were each fitted with seven 16-gauge hypodermic needles (Fig. 1). Randomized drips were emitted from the needles over the exposed upper surface area of the porous media in each column. The manifolds were connected...
to a peristaltic pump using food grade, autoclavable tubing (06402-13, Cole-Parmer, Vernon Hills, IL).

The drip emitters were mounted on top of a vented, foillined acrylic box in which two 254-nm wavelength, ultraviolet (UV) germicidal lamps were housed. The box was attached to the tops of the columns as depicted in Fig. 1. The UV lamps were activated on regular intervals to maintain the sterility of the sand and dripers at the surface of the columns. Auxiliary experiments indicated that the UV light only penetrated the sand a distance equivalent to a few grains or less, so it only sterilized the upper surfaces of the sand in the columns and did not impact biomass growth elsewhere within sand pack.

Effluent samples were collected periodically during the experiments to measure DO using a YSI model 5300 biological oxygen meter with a model 5331 oxygen electrode (Yellow MMS solution was then started on the tops of the columns at 

\[ \frac{d}{dt} \left( \frac{dS_a}{dz} \right) = \frac{d}{dz} \left[ K_z \left( \frac{dh}{dz} + 1 \right) \right] \pm \Omega \]

where \( \phi \) is the porosity, \( S_a \) is the aqueous-phase saturation, \( t \) is time, \( z \) is the vertical coordinate, \( K_z \) is the saturated hydraulic conductivity; \( k_r \) is the relative hydraulic conductivity, \( h \) is the soil water pressure head, and \( \Omega \) is a volumetric flux due to liquid sources or sinks.

A fixed head (Dirichlet) boundary condition was applied to the lower boundary:

\[ h(z, t) = h_0(t); z = L \]

and a prescribed flux (Neumann) boundary was specified for the upper boundary:

\[ -K \frac{dh}{dz} + 1 = q_{so}(z, t); z = 0 \]

where \( h_0 \) and \( q_{so} \) are the prescribed pressure head and water flux, respectively.

Solute and bacterial transport, cell growth, substrate consumption, and gas diffusion were modeled using advection–dispersion reaction equations of the form:

\[ \frac{\partial}{\partial t}\left(R_tC_kC_k\right) = \frac{\partial}{\partial z}\left[ \frac{\partial C_k}{\partial z} \right] - \frac{\partial}{\partial z} (q_{so}C_k) + \lambda_k \]

where \( R_T \) is a dimensionless retardation coefficient, \( \theta \) is the volumetric fluid content, \( C \) is the mass of a particular constitut-
unt per volume of pore fluid, \( D \) is a hydrodynamic dispersion coefficient, \( q \) is the Darcian flux, and \( \Lambda \) represents a reaction-rate source–sink term. Subscript \( k \) refers to different constituents (e.g., \( k = \text{“g”} \) for glucose, \( O_2 \), \( CO_2 \), or \( \text{“m”} \) for microbes), and subscript \( \ell \) refers to the phase (e.g., \( \ell = \text{“w”} \) for water and \( \text{“a”} \) for air).

For aqueous phase solutes, the hydrodynamic dispersion coefficient was defined as
\[
D_{k,w} = \alpha_k |v_w| + D_{k,w}^{\text{eff}}
\]  
where \( \alpha_k \) is the dispersivity, \( v_w \) is the water velocity \((q/\theta)\), \( D_{k,w}^{\text{eff}} \) is the effective molecular diffusion coefficient in water. For \( O_2 \) and \( CO_2 \), which partition between the aqueous and gas phases, dispersion in the gas phase was neglected and effective values of the retardation factor, dispersion coefficient, and Darcian flux were defined in a manner analogous to that used by Šimunek and Suarez (1993):
\[
\begin{align*}
R_l &= 1 + \left( \frac{M_w}{K_h RT} \right) \frac{\theta_a}{\theta_w}, \\
\theta_a D_E &= \theta_a D_{k,w} + \left( \frac{M_w}{K_h RT} \right) \theta_a D_{k,g}^{\text{eff}}, \\
q_E &= q_w + \left( \frac{M_w}{K_h RT} \right) q_a
\end{align*}
\]  
where \( M_w \) is the molecular weight, \( K_h \) is the Henry’s Law constant, \( R \) is the ideal gas constant, and \( T \) is absolute temperature. Effective diffusion coefficients were defined as
\[
D_{k,g}^{\text{eff}} = D_{k,w}^{\text{eff}} \left( \frac{\theta_a}{\theta_w} \right)
\]  
where \( a \) and \( e \) are empirical parameters (Millington and Quirk, 1961; Moldrup et al., 2000). Assuming no gas flow across the lower boundary of the columns, the flux of air into or out of the top of the columns during a time step was calculated from
\[
q_a(0) = -\frac{L}{A} \int_0^t V d\theta_a dz dt
\]  
where \( V \) is the volume of porous media represented by a model grid block, and \( A \) is the cross-sectional area of the grid block normal to the direction of flow. Vertical air fluxes at other locations were estimated in the same manner.

Boundary conditions for Eq. [4] were handled in a manner similar to that described by Šimunek and Suarez (1993). For the upper boundary, first- or third-type boundary conditions were used
\[
\begin{align*}
C_{k,a}(0,t) &= C_{k,a0}, \\
-\theta_a D_E \frac{\partial C_{k,a}}{\partial z} + q_E C_{k,a} &= q_{1b} C_{k,a0}
\end{align*}
\]  
where \( q_{1b} \) is the effective flux, and \( C_{k,a0} \) is the concentration associated with this flux or prescribed at the boundary. For constituents that exist only in the aqueous phase or partition only between the aqueous and solid phases, the third-type boundary condition was used, and \( D_E \) and \( q_E \) did not account for the gas phase. For constituents that partition between the aqueous and gas phases, Eq. [6] through [8] define the effective values of \( R_l, D_E, \) and \( q_E \), which were used in Eq. [4]. The third-type boundary condition given by Eq. [12] was used for multiphase constituents whenever \( q_E > 0 \), and the first-type boundary condition given by Eq. [11] was used when \( q_E \leq 0 \). In both cases, \( C_{k,a0} = C_{k,wa0} \), which corresponds to the equilibrium concentration in the aqueous phase determined by the atmospheric concentration of the gas using Henry’s Law. For the lower boundary, a continuous concentration boundary condition was specified:
\[
\frac{\partial C_{k,a}(L,t)}{\partial z} = 0
\]  

The coefficients of molecular diffusion for \( O_2 \) and \( CO_2 \) in air were estimated from (Jaynes and Rogowski, 1983):
\[
D_{\text{O}_2,a}^{\text{mol}} = \frac{D_{\text{O}_2,C_0,N_2}}{X_{\text{O}_2,C_0,N_2} + D_{\text{O}_2,C_0,N_2} + rD_{\text{O}_2,C_0,N_2}}
\]  
\[
D_{\text{CO}_2,a}^{\text{mol}} = \frac{D_{\text{CO}_2,C_0,N_2}}{X_{\text{CO}_2,C_0,N_2} + D_{\text{CO}_2,C_0,N_2} + (D_{\text{CO}_2,C_0,N_2} + r)}
\]  

Reactions were assumed to occur only in the aqueous phase and were represented using the following equations:
\[
\begin{align*}
\Lambda_{g,w} &= -\mu_{Y,g,w}(\theta_m C_{m,w} + \rho_0 C_{m,s1}), \\
\Lambda_{g,a} &= -\mu_{Y,g,a}(\theta_m C_{m,a} + \rho_0 C_{m,s1}), \\
\Lambda_{m,w} &= \mu_{Y,m,w} C_{m,w} - k_1 \theta_a C_{m,a} + k_2 \rho_0 C_{m,s1}
\end{align*}
\]  

where \( \mu \) is the specific growth rate of the bacterium; \( Y_{g,w} \) and \( Y_{g,a} \) are yield coefficients representing the mass of electron donor or substrate (e.g., glucose) consumed and the mass of terminal electron acceptor (e.g., \( O_2 \)) consumed, respectively, per mass of cells generated; \( k_1 \) and \( k_2 \) are first-order reversible attachment and detachment coefficients, respectively; and \( \rho_0 \) is the bulk density. The terms \( C_{m,w} \) and \( C_{m,a} \) represent the mass of cells in the aqueous phase per volume of pore liquid, and the mass of cells reversibly attached to solids per mass of porous media, respectively.

The specific growth rate was represented by a multiplicative Monod-type kinetics model (McGee et al., 1970):
\[
\mu = \mu_{\text{max}} \left( \frac{C_{o,w}}{K_o + C_{o,w}} \right) \left( \frac{C_{g,w}}{K_g + C_{g,w}} \right)
\]  

where \( \mu_{\text{max}} \) is the maximum specific growth rate \((h^{-1})\), and \( K_o \) and \( K_g \) are half-saturation (or half-saturation) coefficients \((\text{mg} \ L^{-1})\) for \( O_2 \) and the glucose, respectively.

The mass balance equation for the attached biomass was
\[
\frac{\partial}{\partial t} (\rho_s C_{m,s1}) = \mu \rho_0 C_{m,s1} + k_1 \theta_a C_{m,w} - k_2 \rho_0 C_{m,s1}
\]  

where the attachment coefficient, \( k_1 \), was estimated from (Tien et al., 1979)
\[
k_1 = \frac{3}{2} \frac{d_i}{\theta_a} \frac{(1 - \theta_a)}{d_i} \eta \alpha_c
\]  

where \( d_i \) is the median grain diameter of the porous medium (or collector), and \( \eta \) and \( \alpha_c \) are the so-called collector and collision (or sticking) efficiencies (Logan et al., 1995; Deshpande and Shonnard, 1999). The \( \eta \) parameter was calculated using the particle-filtration model of Rajagopalan and Tien.
(1976). Cell attachment–detachment kinetics were assumed to be related to local environmental conditions, including the presence of growth substrate. The detachment coefficient, $k_d$, was represented by

$$k_d = f \exp(-tg_c)$$  \[22\]

where $f$ and $g_c$ are empirical parameters, and $t$ was taken as the time since any location in the modeled domain was first exposed to a minimum threshold concentration of growth substrate (1 $\mu$g L$^{-1}$ glucose). Although the form of Eq. [22] is somewhat arbitrary, it was found to adequately reproduce the observed behavior of the effluent biomass. Time-dependent attachment–detachment kinetics have been applied previously to model virus and protein adsorption (Lee et al., 1999). Alternative approaches for modeling cell attachment and detachment have been described by Ginn (1999), Ginn et al. (2002), and Johnson et al. (1995).


**Parameter Estimation**

The saturated hydraulic conductivity, $K_s$, of the 40/50 Accusand was determined using the falling head method (Klute, 1986). The water retention characteristics of the clean sand were determined by wet-packing the column, allowing it to drain to a hydrostatic condition, and then sampling each column segment to determine gravimetric water contents. These volumetric water contents were then used to calculate gravimetric water contents using the overall bulk density of the sand pack. The volumetric water contents were paired with pressure heads for each depth, which were taken as the negative value of the elevation above the location of zero pressure in the hydrostatic column.

The water retention characteristics of the sand were represented using the model of Brooks and Corey (1964). The relative permeability or unsaturated hydraulic conductivity of the sand was represented using the model of Burdine (1953). The hydraulic properties of the attached biomass phase were represented using the models of van Genuchten (1980) and Mualem (1976) with the parameters of a clay soil. The sand and assumed biomass properties were combined using the composite media model described by Rockhold et al. (2002). The hydraulic properties for the sand and biomass are depicted in Fig. 2 and parameter values are given in Table 1.

The parameters in Eq. [19] were estimated from batch growth experiments. The stoichiometry of the biologically mediated redox reaction was estimated using the energetics model of McCarty (1975). This model predicts the following stoichiometry, assuming 60% efficiency in the conversion of glucose to cell biomass (C$_6$H$_{12}$O$_6$N), with NH$_4^+$ as the source of N, and with O$_2$ as the terminal electron acceptor

\[ 1.15 \text{C}_6\text{H}_2\text{O}_6 + 1.9 \text{O}_2 + \text{NH}_4^+ + \text{HCO}_3^- \rightarrow \text{C}_6\text{H}_4\text{O}_2\text{N} + 2.9 \text{CO}_2 + 5.9\text{H}_2\text{O} \]  \[23\]

HK44 is a facultative aerobe capable of denitrification. Therefore a source of NH$_4^+$ was included in the MMS growth media, rather than NO$_3^-$, to prevent HK44 from using NO$_3^-$ as an alternative electron acceptor if anoxic conditions developed. Equation [23] indicates that 1.15 mol of glucose and 1.9 mol of O$_2$ are required to produce 1 mol of bacterial cells. This translates into $Y_m = 1.83$ (mg glucose mg$^{-1}$ cells), and $Y_m = 0.54$ (mg O$_2$ mg$^{-1}$ cells). Independent batch experiments were also conducted to estimate yield coefficients. The experimentally determined values were variable, but similar to the estimates given by Eq. [23]. Equation [23] was also used to estimate the respiration coefficient, $r_{CO_2} = 1.53$ (mol CO$_2$ produced mol$^{-1}$ O$_2$ consumed). Table 1 summarizes the parameters that were used for the base case model simulation.

The initial water flow in the columns was specified as a uniform pressure head, $h(x,0) = 0$ cm (fully saturated). The upper boundary condition was a constant flux, $q_s(0,t) = 7$ cm h$^{-1}$. The lower boundary condition was specified as a held pressure head that was adjusted in steps to mimic the slow release of water from the columns during the early stages of the experiment, followed by a fixed pressure head of 1 cm at later times. The pressure head steps were: $h(0 < t < 0.01) = 43$ cm, $h(0.01 < t < 0.02) = 39$ cm, $h(0 < t < 0.03) = 35$ cm, $h(0.03 < t < 0.04) = 30$ cm, $h(0.04 < t < 0.05) = 25$ cm, $h(0.05 < t < 0.06) = 20$ cm, $h(0.06 < t < 0.07) = 15$ cm, $h(0.07 < t < 0.08) = 10$ cm, $h(0.08 < t < 0.1) = 5$ cm, $h(0.01 < t < 0.1) = 1$ cm.

The initial aqueous-phase glucose, O$_2$, and CO$_2$ concentrations were specified as 0.0, 9.2, and 0.01 mg L$^{-1}$, respectively. These initial O$_2$ and CO$_2$ concentrations correspond to equilibrium with gas phase concentrations in the atmosphere as determined using Henry’s Law. The upper boundary condition for glucose was a constant concentration of 250 mg L$^{-1}$. Initial gas-phase concentrations and upper boundary conditions for O$_2$ and CO$_2$ were specified to correspond with equilibrium conditions with the atmosphere. The initial aqueous-phase cell concentration was specified as 80.4 mg L$^{-1}$ ($\approx 3 \times 10^8$ CFU mL$^{-1}$) everywhere, except for the top node in the model grid, where a fixed concentration of 0.0 mg L$^{-1}$ was specified to account for the presence of the UV germicidal lamp, which was turned on at regular intervals during the experiment to prevent cell growth on the surface of the sand. It was further assumed that no bacteria were initially attached to the sand anywhere in the columns.

**RESULTS AND DISCUSSION**

Increases in pressure heads and water contents were observed during the experiments. One possible mechanism that could cause this behavior is lowering of surface tension due to adsorption of cells and/or biosurfactants at air–water interfaces. Lowering of surface tension would tend to cause desaturation of the porous media, rather than the increases in water content that were observed. Therefore this mechanism may initially seem counterintuitive. However, localized desaturation would reduce the unsaturated hydraulic conductivity of the porous media, which would effectively increase the resistance to flow, possibly leading to increases in upstream saturations. Independent measurements of surface tension for aqueous solutions of the MMS medium with different concentrations of stationary-phase HK44 cells indicated that surface tension lowering due to sorption of stationary-phase cells alone was not significant ($<1$ mN m$^{-1}$) for cell concentrations up to approximately $10^6$ CFU mL$^{-1}$ (Rockhold et al., 2002). However, surface tension was not measured on suspensions of cells that were meta-
bolically active, or under conditions of O\textsubscript{2} stress, or for cell densities $>10^9$ CFU mL\textsuperscript{-1}.

Studies reported by Déziel et al. (1996) and Kosaric (1993) suggest that metabolically active cells may generate surface-active agents and possibly change their surface character to become more hydrophobic. The generation of surface-active compounds (e.g., excess fatty acids) might have resulted in surface tension lowering in excess of what was observed for stationary-phase cells alone. The anoxic environment that developed during high substrate loading conditions of the experiment might also have promoted the migration of the motile bacteria to air–water interfaces via chemotaxis. Sorption of bacterial cells on the surfaces of mineral grains may also change the wettability of a porous medium, which could result in increases in apparent contact angles (Abelson et al., 1983).

Decreases in surface tension and increases in apparent contact angle are both a function of changes in interfacial energies and therefore tend to manifest similar effects. These types of changes were accounted for by similar-media scaling (Rockhold et al., 2002). The apparent surface tension lowering effect was represented by scaling the capillary pressure–saturation curves using scaling factors calculated as

$$\frac{\sigma}{\sigma_0} = 1 - \frac{1}{A C_{m,w}^B}$$

where $\sigma_0$ is the surface tension of the cell-free MMS solution (74.2 mN m\textsuperscript{-1}), $A$ and $B$ are empirical parameters, and $C_{m,w}$ is the aqueous-phase cell concentration. Equation [24] is plotted in Fig. 3 using the $A$ and $B$ parameters from Table 1.

Equation [24] was required to approximately reproduce the observed changes in apparent saturations and pressure heads in the columns. It should be noted, however, that capillary pressure represents the difference between the nonwetting (air) and wetting fluid (water) pressures. Therefore if the total gas pressure increases, this could also result in an apparent surface tension lowering effect. The surface tension lowering function serves to capture the effects of possible changes in fluid-media properties, including adsorption of surface-active compounds at gas–liquid interfaces, as well as possible increases in total gas pressure. Changes in total gas pressure due to CO\textsubscript{2} production could be modeled more mechanistically, however, by solving fully coupled equa-

### Table 1. Parameters used for model simulations.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Base case value (units)</th>
<th>Source or reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>$K_{\text{mat}}$</td>
<td>317 (cm h\textsuperscript{-1})</td>
<td>independent measurement</td>
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<tr>
<td>$\theta_0$</td>
<td>0.37 (cm\textsuperscript{2} cm\textsuperscript{-3})</td>
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<tr>
<td>$\theta_{\text{m,w}}$</td>
<td>0.028 (cm\textsuperscript{2} cm\textsuperscript{-3})</td>
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<td>$\lambda_{\text{m}}$</td>
<td>21.6 (cm)</td>
<td>&quot;</td>
</tr>
<tr>
<td>$K_{\text{col}}$</td>
<td>5.84 &quot;</td>
<td></td>
</tr>
<tr>
<td>$K_{\text{biomass}}$</td>
<td>1.08 (cm h\textsuperscript{-1})</td>
<td>clay soil (Leij et al., 1999)</td>
</tr>
<tr>
<td>$\theta_{\text{biomass}}$</td>
<td>0.51 (cm\textsuperscript{2} cm\textsuperscript{-3})</td>
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</tr>
<tr>
<td>$\phi_{\text{biomass}}$</td>
<td>0.102 (cm\textsuperscript{3} cm\textsuperscript{-3})</td>
<td>&quot;</td>
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<tr>
<td>$\sigma_{\text{biomass}}$</td>
<td>0.021 (cm\textsuperscript{-1})</td>
<td>&quot;</td>
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<td>$\rho_{\text{biomass}}$</td>
<td>1.2 &quot;</td>
<td></td>
</tr>
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<td>$\rho_{\text{composite media}}$</td>
<td>3.17 Clement et al. (1996)</td>
<td></td>
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<td>$\rho_{\text{s}}$</td>
<td>1730 (mg cm\textsuperscript{-3})</td>
<td>independent measurement</td>
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<td>$\rho_{\text{b}}$</td>
<td>1100 (mg cm\textsuperscript{-3})</td>
<td>Bouwer and Rittmann (1992)</td>
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<td>independent measurement</td>
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<td>Barton and Ford (1995)</td>
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<td>Marrero and Mason (1972)</td>
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<tr>
<td>$D_{\text{ch}_{\text{N}},\text{o}}$</td>
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<td>$D_{\text{CO}_{2},\text{o}}$</td>
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<tr>
<td>$K_{\text{H}_{2}\text{CO}_3}$</td>
<td>29.5 ([L atm] (mol K\textsuperscript{-1})\textsuperscript{-1})</td>
<td>&quot;</td>
</tr>
<tr>
<td>$\alpha$</td>
<td>2.42</td>
<td>fit to experimental data</td>
</tr>
<tr>
<td>$A$</td>
<td>60</td>
<td>fit to experimental data</td>
</tr>
<tr>
<td>$B$</td>
<td>0.54</td>
<td>fit to experimental data</td>
</tr>
<tr>
<td>$\sigma_0$</td>
<td>74.2 (mN m\textsuperscript{-1})</td>
<td>independent measurement</td>
</tr>
<tr>
<td>$\mu_{\text{biomass}}$</td>
<td>0.496 (h\textsuperscript{-1})</td>
<td>&quot;</td>
</tr>
<tr>
<td>$K_{\text{a}}$</td>
<td>10.7 (mg L\textsuperscript{-1})</td>
<td>&quot;</td>
</tr>
<tr>
<td>$K_{\text{r}}$</td>
<td>1.5 (mg L\textsuperscript{-1})</td>
<td>&quot;</td>
</tr>
<tr>
<td>$Y_{\text{a}}$</td>
<td>1.83 (mg mg\textsuperscript{-1})</td>
<td>fit to experimental data</td>
</tr>
<tr>
<td>$Y_{\text{r}}$</td>
<td>0.54 (mg mg\textsuperscript{-1})</td>
<td>Eq. [23]</td>
</tr>
<tr>
<td>$Y_{\text{v}}$</td>
<td>1.13 (mg mg\textsuperscript{-1})</td>
<td>&quot;</td>
</tr>
<tr>
<td>$R$</td>
<td>1.53 (mol CO\textsubscript{2} mol\textsuperscript{-1} O\textsubscript{2})</td>
<td>&quot;</td>
</tr>
<tr>
<td>$f$</td>
<td>0.29 (h\textsuperscript{-1})</td>
<td>fit to experimental data</td>
</tr>
<tr>
<td>$g$</td>
<td>0.014 (h\textsuperscript{-1})</td>
<td>&quot;</td>
</tr>
</tbody>
</table>
The simulation results match the observed effluent glucose and biomass concentration data reasonably well. However, the DO data indicate relatively constant concentrations of approximately 1 mg L\(^{-1}\) in the effluent after the initial drainage period, whereas the simulation results indicate that the concentration of DO in the effluent at these times was near zero. Differences between observed and simulated effluent DO concentrations may be due to (i) partial re-oxygenation of the effluent DO samples as they were transferred from the sampling syringes on the columns to the measurement vessel, (ii) diminished HK44 respiration rates at low DO concentrations that were not accounted for in the model, and/or (iii) gas–liquid mass transfer limitations associated with adsorption of cells and/or biosurfactants at gas–liquid interfaces.

The equilibrium concentration of O\(_2\) in deionized water at atmospheric pressure with a partial O\(_2\) gas pressure of 0.021 MPa (0.21 atm) is 9.2 mg L\(^{-1}\) (Lide, 1996). The stoichiometry indicated by Eq. [23] suggests that for this DO concentration, only about 31 of the 250 mg L\(^{-1}\) of glucose in the influent would be consumed, leaving a concentration of 219 mg L\(^{-1}\) glucose in the column effluent. Simulations that neglected gas diffusion verified this...
calculation. However, gas diffusion results in O\textsubscript{2} being replenished at a much faster rate than is possible by advection and diffusion in the aqueous phase alone, thus allowing virtually all of the glucose to be utilized. This result demonstrates the importance of gas-phase diffusion to soil respiration.

The gradual decrease in effluent biomass concentrations that was observed after about 60 h (Fig. 4) required the use of a time-dependent detachment rate coefficient (Eq. [22]) to reproduce this behavior in the model simulations. However, this effect might also have been achieved by accounting for biomass decay or cell death. Decreases in effluent biomass concentrations might also have been due, at least in part, to decreases in the effective growth rate of the bacteria, possibly resulting from gas–liquid or liquid–cell mass transfer limitations that developed with time (Rockhold et al., 2004). For example, Bailey and Ollis (1986) noted that for a variety of sparingly soluble gases, surfactant adsorption at gas–liquid interfaces resulted in an average reduction in the interphase mass transfer coefficient of 60%. The observed and simulated results shown in Fig. 4 deviate somewhat between times of about 5 and 50 h, with the observed effluent data exhibiting lower concentrations than the simulation results between 5 and 30 h, and higher concentrations than the simulation results between 40 and 50 h. These differences may be attributable to other phenomena such as irreversible sorption of bacteria on Fe oxide coatings on the sand and/or on air–water interfaces.

Figure 4 also shows the observed and simulated final sand-associated biomass concentration distributions in the columns. The apparently abrupt increase in attached cell concentrations indicated by the simulation results at a depth of approximately 30 cm reflects the position of the capillary fringe. Below a depth of about 32 cm, the sand is fully water saturated, with a volumetric water content, $\theta_w = 0.37$. Water contents decrease rapidly above this depth, and at the 27-cm depth $\theta_w = 0.15$. In the upper, unsaturated part of the columns, ample O\textsubscript{2} supply, which is replenished by gas-phase diffusion, allows for more cell growth relative to the lower saturated parts of the columns. The attachment coefficient, $k_1$, also increases as water content decreases, as indicated by Eq. [21].

Figure 5 shows time histories of observed and simulated values of volumetric water content and pressure head at selected locations. The observed changes in water contents and pressures could only be reproduced approximately in the model simulations by using fluid-media scaling to account for an apparent surface-tension lowering effect. As shown in Fig. 6, significant increases in apparent volumetric water contents were only observed at the uppermost TDR measurement locations, approxi-

![Image](image_url)
Water content measurement by TDR is a function of the dielectric permittivities of the materials in which the probes (or waveguides) are embedded (e.g., water, air, sand, and biomass). The energy of the electromagnetic (EM) pulse that is propagated and reflected back during a measurement is most highly concentrated in the immediate vicinity of the waveguides, and decreases rapidly with distance away from them (Knight, 1992). At higher water contents, the energy of the EM pulse would not propagate as far away from the waveguides because of the higher dielectric permittivity of water (≈80) relative to air (≈1). It is possible that the uppermost TDR measurements were influenced by the accumulation of biomass. Bacterial cells are typically on the order of 70 to 90% water by volume (Madigan et al., 1997). Therefore they might be expected to appear similar to water in terms of an apparent water content measured by TDR. However, the molecular structure of bacterial cells is certainly different than water and would have different relaxation frequencies. The TDR measurements may be susceptible to bias of this type due to the high bandwidth of the Tektronix Model 11801 oscilloscope. The volumes sampled by the TDR measurements at the 1.6-cm depth may extend up to the surface of the sand in the columns and become progressively more biased with time because of the higher biomass concentrations that develop near the surface. The possible influence of bacterial cell accumulation on TDR measurements was not investigated or accounted for in the TDR probe calibrations.

Figure 5 also shows time histories of observed and simulated pressure head values, respectively, at selected measurement depths. Two of the tensiometer segments in Column A developed air leaks that rendered the data from those locations useless. Unlike the TDR data, which showed significant increases in apparent volumetric water contents only near the surface, the transducer data indicate significant increases in pressure heads at all depths. A steady water flux was established on the surfaces of the columns. Therefore differences between measured values of pressure heads at different depths that should all be under a unit hydraulic gradient may be indicative of nonuniformities in packing and/or bacterial-induced changes in the hydraulic properties of the sand.

The transducer data for Column B show sudden increases in the rates of change of pressure heads at a time of approximately 145 h. These increases in pressure head correspond with the start of ponding of the influent on the surface of Column B that occurred at this time, as noted in Fig. 5. At this point, the hydraulic conductivity of the sand in Column B was reduced from its original saturated value of approximately 317 cm h⁻¹ to <7 cm h⁻¹, which represents a 45-fold decrease. Ponded water at the surface effectively reduces the air permeability to zero, preventing gas from either entering or exiting the column. The sudden rate of increase in pressure heads after 145 h is due to compression of gas trapped below the ponded water while infiltration continued, as well as possible gas pressure buildup due to the continued production of CO₂ by bacteria within the column. After 145 h, the single-phase Richards equation is clearly no longer valid for modeling these systems.

The calculated partial pressures of O₂ and CO₂ in the columns at the top of the capillary fringe (≈30 cm depth) just before destructive sampling (168 h) were approximately 0.0001 and 0.035 MPa (0.001 and 0.35 atm), respectively. In order for atmospheric pressure to be maintained in this system, the increase in the partial pressure of CO₂ above 0.021 MPa would require the partial pressure of N₂ to decrease from about 0.079 to 0.065 MPa. If N₂ is more or less stagnant, nonequimolar respiration \((r > 1)\) would tend to increase the total gas pressure in the system as a result of excess production of CO₂. If the upper boundary of the system is maintained at atmospheric pressure, even very slight increases in total gas pressure within the columns would lead to some advective movement of gases out of the system, rather than just the simple Fickian diffusion process that was represented in the model (Leffelaar, 1988; Freijer and Leffelaar, 1996). Total gas pressures would tend to adjust so that they remain close to atmospheric pressure, as long as the gas phase is continuous throughout the porous media, and resistance to air flow is negligible. Additional considerations for modeling gas transport in unsaturated porous media are discussed by Thorstenson and Pollock (1989).

Figure 6 shows the measured water retention characteristics and fitted Brooks–Corey model representing...
the primary drainage curve for the clean 40/50 Accu-sand, and pressure–saturation values from various depths in Column B that were calculated from the transducer and TDR data during the experiment. Also shown in Fig. 6 is a hypothetical wetting curve resulting from simple scaling of the drainage curve, using an assumed contact angle of 30°. It is not known whether the apparent changes in water content and pressure heads that were observed in these experiments resulted primarily from (i) clogging of pore throats by cell aggregates or biofilms, (ii) CO2 gas generation and occlusion, (iii) production of biosurfactants (e.g., excess fatty acid production due to the high substrate loading conditions) with concomitant lowering of gas–liquid interfacial tension, (iv) changes in the wettability of the sand as it became coated by biofilms, or (v) a combination of these factors. A combination of these factors is most likely. Regardless of the mechanism(s), it is clear from these experiments that biomass-induced changes in the hydraulic properties of variably saturated porous media can be significant. Such changes are not typically considered when modeling flow and reactive transport processes in soils.

SUMMARY AND CONCLUSIONS

An experimental and numerical investigation was conducted to study interactions between microbial dynamics and transport processes in sand-packed columns under initially steady, variably saturated flow conditions. Glucose was used as the sole C source for growth of Pseudomonas fluorescens HK44. Increases in pressure heads were observed at all measured depths during the experiment, but significant increases in apparent volumetric water contents only occurred within the upper 5 cm of the columns, corresponding to the areas of highest attached biomass concentrations. The hydraulic conductivity of the sand was reduced by a factor of 45 during the 1-wk experiment. To the best of our knowledge, this is one of the first experiments to document bacterially induced changes in the hydraulic properties of unsaturated porous media during active cell growth and transport.

The experiment was simulated using the single-phase Richards equation to model water flow and advection–dispersion reaction reactions to model solute and bacterial transport, cell growth and accumulation, glucose and O2 consumption, and gas diffusion and exchange. Reasonably good matches were obtained between observed and simulated effluent glucose and biomass data and final sand-associated (attached) biomass concentration distributions using first-order reversible attachment–detachment kinetics with attachment coefficients based on particle-filtration theory, and time-dependent detachment rate coefficients. The observed changes in water contents and pressure heads were reproduced approximately using fluid-media scaling to account for an apparent surface-tension lowering effect.

During the experiment, ponded water developed on the surface of one of the columns after about 6 d. Pressure head readings increased rapidly after this time, while water contents remained relatively constant, indicating compression of entrapped gas and possible buildup of gas pressure due to continued production of CO2 by bacteria. The single-phase Richards equation was no longer valid for modeling the experiment after this point because the air phase was discontinuous.

These results illustrate the potential importance of using fully coupled multifluid flow equations, rather than the single-phase Richards equation, and multicomponent reactive transport equations to model coupled biogeochemical reactions and transport in soils.

Considerable uncertainties remain as to the exact mechanism(s) responsible for the changes in hydraulic properties that were observed during this study, suggesting that further experimental work is warranted. Using additional instrumentation such as microelectrodes, (De-Beer and Schramm, 1999; Lewandowski et al., 1999) and in-line gas sampling, might allow a more complete understanding of the mechanisms and process interactions.

REFERENCES


Estrella, M.R., M.L. Brusseau, R.S. Maier, I.L. Pepper, P.J. Wierenga,


