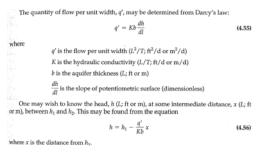
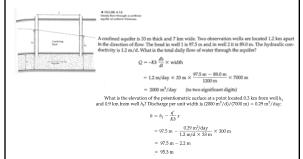


Confined Aquifer Steady Flow Calculations



Confined Aquifer Steady Flow Calculations



Unconfined Aquifer Steady Flow Calculations



◆ FIGURE 4.17

Steady flow through an unconfined aquifer resting on a horizontal impensions surface.

This problem was solved by Dupuit (1863), and his assumptions are known as the Dupuit assumptions. The assumptions are that (1) the hydraulic gradient is equal to the slope of the water table and (2) for small water-table gradients, the streamlines are horizontal and the equipotential lines are vertical. Solutions based on these assumptions have proved to be useful in many practical problems. However, the Dupuit assumptions do not allow for a seepage face above the outflow side.

From Dany's law,

$$q' = -Kh \frac{dh}{dx} ag{4.57}$$

where h is the saturated thickness of the aquifer. At x=0, $h=h_{1i}$ at $x=L_{i}$, $h=h_{2}$. Equation 4.57 may be set up for integration with the boundary conditions:

$$\int_{0}^{L} q' dx = -K \int_{b_{j}}^{b_{j}} h dh$$

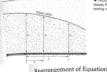
Integration of the preceding yields

$$q'x\Big|_0^L = -K\frac{h^2}{2}\Big|_{h_1}^{h_2}$$

Substitution of the boundary conditions for x and h yields

$$q'L = -K\left(\frac{h_2^2}{2} - \frac{h_1^2}{2}\right)$$

Unconfined Aquifer Steady Flow Calculations **ROBE 4.17 Steady Str. House, or uncortend aquifers strenge on a functional aquifers



Rearrangement of Equation 4.58 yields the Dupuit equation:

$$q' = \frac{1}{2}K\left(\frac{h_1^2 - h_2^2}{L}\right)$$

where

q' is the flow per unit width $(L^2/T; \mathrm{ft}^2/\mathrm{d} \ \mathrm{or} \ \mathrm{m}^2/\mathrm{d})$

K is the hydraulic conductivity (L/T; ft/d or m/d)

 h_1 is the head at the origin (L; ft or m)

 h_2 is the head at L(L; ft or m)

L is the flow length (L; ft or m)

▲ FIGURE 4.18
Control volume for flow through a prism of an unconfined aquifer with the bottom resting on a horizontal impervious surface and the top coinciding view water table.



(4.59)

If we consider a small prism of the unconfined aquifer, it will have the shape of Figure 4.18. On one side it is h units high and slopes in the x-direction. Given the Duputl assumptions, there is no flown in the x-direction. The flow in the x-direction. The flow in the x-direction of the prism is x-direction of the prism is x-direction through the left face of the prism is

$$q'_{x}dy = -K\left(h\frac{\partial h}{\partial x}\right)_{x}dy$$
 (4.6)

where dy is the width of the face of the prism. The discharge through the right face, q'_{x+dx} is

$$q'_{x+dx}dy = -K\left(h\frac{\partial \dot{\theta}}{\partial x}\right)_{x+dx}dy$$
 (4.61)

Note that $\left(h,\frac{\partial h}{\partial x}\right)$ has different values at each face. The change in flow rate in the x-direction between the two faces is given by

$$(q'_{s+dx} - q'_s)dy = -K\frac{\partial}{\partial x}\left(h\frac{\partial h}{\partial x}\right)dx dy$$
 (4.62)

Through a similar process, it can be shown that the change in the flow rate in the y-direction is

$$(q'_{y+dy} - q'_{y})dx = -K \frac{\partial}{\partial y} \left(h \frac{\partial h}{\partial y}\right) dy dx$$

(4.63)



For steady flow, any change in flow through the prism must be equal to a gain or loss of water across the water table. This could be infiltration or evapotranspiration. The net addition or loss is at a rate of w, and the volume change within the initial volume is w dx dy where dx dy is the area of the surface. If w represents evapotranspiration, it will have a negative value. As the change in flow is equal to the new addition,

change in flow is equal to the new addition,
$$-K\frac{\partial}{\partial x}\left(h\frac{\partial h}{\partial x}\right)dx\,dy - K\frac{\partial}{\partial y}\left(h\frac{\partial h}{\partial y}\right)dy\,dx = w\,dx\,dy \tag{4.64}$$

We can simplify Equation 4.64 by dropping out $dx\,dy$ and combining the differentials:

$$-K\left(\frac{\partial^2 h^2}{\partial x^2} + \frac{\partial^2 h^2}{\partial y^2}\right) = 2w \tag{4.65}$$

If w=0, then Equation 4.65 reduces to a form of the Laplace equation:

$$\frac{\partial^2 h^2}{\partial x^2} + \frac{\partial^2 h^2}{\partial y^2} = 0 {4.66}$$



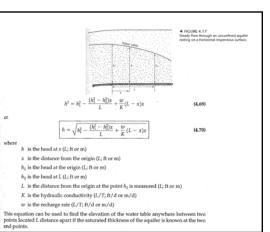
If flow is in only one direction and we align the x-axis parallel to the flow, then there is no flow in the y-direction, and Equation 4.65 becomes

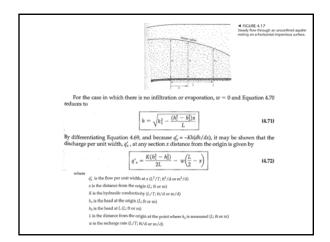
$$\frac{d^2(h^2)}{dx^2} = -\frac{2w}{K} {4.62}$$

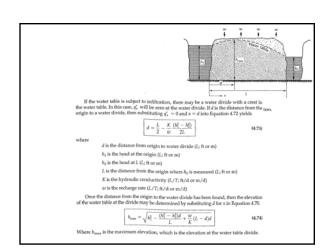
Integration of this equation yields the expression

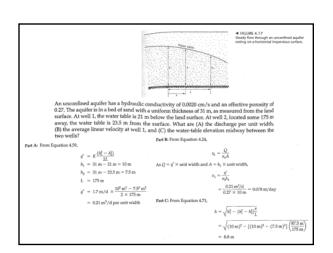
$$h^2 = -\frac{wx^2}{K} + c_1x + c_2 (4.68)$$

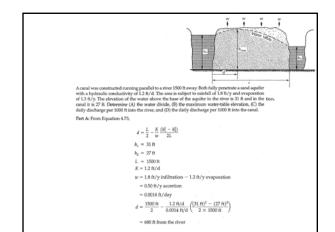
where c_1 and c_2 are constants of integration. The following boundary conditions can be applied: at x=0, $h=h_1$; at x=L, $h=h_2$ (Figure 4.19). By substituting these into Equation 4.68, the constants of integration can be evaluated with the following result:

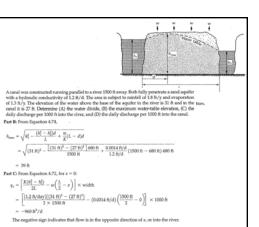


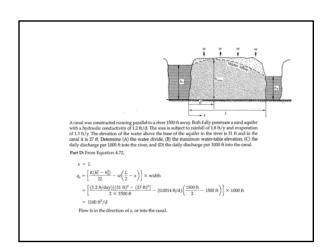












= -960 ft³/d